

**TEKST NR 160**

**1988**

**Loop Groups  
and  
Instantons in Dimension Two**

by

Jens Gravesen

**TEKSTER fra**

**IMFUFA**

**ROSKILDE UNIVERSITETSCENTER**

INSTITUT FOR STUDIET AF MATEMATIK OG FYSIK SAMT DERES  
FUNKTIONER I UNDERVISNING, FORSKNING OG ANVENDELSER

IMFUFA, Roskilde University Center, Postbox 260, DK-4000 Roskilde,  
DENMARK

## Loop Groups and Instantons in Dimension Two

Jens Gravesen

IMFUFA tekst no.160/May 1, 1988

7 pages

ISSN 0106-6242

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**Abstract.** Radial trivialization from the poles gives a map from the space of equivalence classes of connections in a  $G$ -bundle over  $S^2$  to the space of loops  $\Omega G$ . The map is a homotopy equivalence, and if the equivalence class  $[\omega]$  is mapped to the loop  $\gamma$ , then  $\gamma$  together with a map  $f: S^2 \rightarrow \mathfrak{g}$  into the Lie algebra give a complete description of  $[\omega]$ . The Yang-Mills functional taken at  $[\omega]$  corresponds to the sum of the energy of  $\gamma$  and a certain norm of  $f$ . In particular, the moduli space of instantons is the same as the space of homomorphisms  $S^1 \rightarrow G$ .

There is a similar description of connections in a  $G$ -bundle over an arbitrary Riemann surface, but so far not of the Yang-Mills functional and the moduli space.

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**Keywords.** Loop groups, Yang-Mills equations, Riemann surfaces.  
**1980 Mathematics subject classifications:** 32G13

## 1. Instantons on the Sphere

Consider the two dimensional unit sphere

$$S^2 = \{(x, y, z) \in \mathbb{R}^3 \mid x^2 + y^2 + z^2 = 1\}.$$

Denote  $(0, 0, 1)$  by  $\infty$  (the North Pole),  $(0, 0, -1)$  by  $0$  (the South Pole), the complement of  $0$  by  $U_\infty$ , the complement of  $\infty$  by  $U_0$ , the part of  $S^2$  where  $z > 0$  by  $D_\infty$  (the northern hemisphere) and the part of  $S^2$  where  $z < 0$  by  $D_0$  (the southern hemisphere). Finally let  $\varphi \in [0, \pi]$  and  $\theta \in [0, 2\pi]$  be spherical coordinates on  $S^2$ .

Let  $G$  be a compact Lie group with Lie algebra  $\mathfrak{g}$  and let  $P$  be a principal  $G$ -bundle over  $S^2$  with a connection  $\omega$ . Choose a base point  $p_\infty$  in the fiber over  $\infty$ .

By lifting the curves  $t \mapsto (t, \theta)$  for  $\theta \in [0, 2\pi]$  to horizontal curves in  $P$  starting at  $p_\infty$ , we get a section  $\sigma_\infty$  in  $P$  over  $U_\infty$ , and by lifting the curves  $t \mapsto (\pi - t, \theta)$  to horizontal curves in  $P$  starting at some point  $p_0$  in the fiber over  $0$ , we get a section  $\sigma_0$  in  $P$  over  $U_0$ . We say that  $\sigma_0$  and  $\sigma_\infty$  are obtained by *radial trivialization* from the poles. Henceforth  $j$  denotes either  $0$  or  $\infty$ .

(1.1) LEMMA. *The sections  $\sigma_0$  and  $\sigma_\infty$  in  $P$  are smooth.*

PROOF: Use the connection on  $P$  to combine the metric on  $S^2$  and the biinvariant metric on  $G$  to give a metric on  $P$ . Obviously a horizontal lift of a geodesic in  $S^2$  gives a geodesic in  $P$ . Hence  $\sigma_j = \exp_P \circ \pi_*^{-1} \circ \exp_{S^2}^{-1}$ , where  $\pi_*$  is the projection from the horizontal space in  $P$  to  $TS^2$  and  $\exp_M$  is the exponential map from  $TM$  to  $M$ .  $\square$

The transition function  $\gamma: U_0 \cap U_\infty \rightarrow G$  is defined by

$$(1.2) \quad \sigma_0 = \sigma_\infty \gamma.$$

As the curve  $t \mapsto \sigma_j(t, \theta)$  is horizontal,  $t \mapsto \gamma(t, \theta)$  is constant for all  $\theta$ , hence  $\gamma$  can be regarded as a smooth map  $\gamma: S^1 \rightarrow G$ , i.e. as an element of  $LG$ .

As  $p_0$  is arbitrary,  $\sigma_0$  and hence  $\gamma$  is only determined up to multiplication by an element of  $G$ , but we have a well-defined element of  $LG/G = \Omega G$ . Alternatively we can put  $p_0 = \lim_{t \rightarrow \pi} \sigma_\infty(t, 0)$ .

We put

$$(1.3) \quad \omega_j = \sigma_j^*(\omega) \in \Omega^1(E_{\mathfrak{g}}|_{U_j}) \cong C^\infty(\wedge^1(U_j, \mathfrak{g}))$$

where  $E_{\mathfrak{g}} = P \times_G \mathfrak{g}$ .

(1.4) LEMMA. *The loop  $\gamma$  and the forms  $\omega_0, \omega_\infty$  are gauge invariant.*

PROOF: Let  $g$  be a based gauge transformation, i.e. an automorphism  $g: P \rightarrow P$  with  $g(p_\infty) = p_\infty$ . Then  $\tilde{\sigma}_j = g^{-1} \circ \sigma_j$  is the section induced by the gauge transformed connection  $\tilde{\omega} = g^*(\omega)$ . We can write  $\tilde{\sigma}_j = \sigma_j g_j$ , then  $g_0 = \gamma^{-1} g_\infty \gamma$ ,

$$\tilde{\sigma}_0 = \sigma_0 g_0 = (\sigma_\infty \gamma)(\gamma^{-1} g_\infty \gamma) = \sigma_\infty g_\infty \gamma = \tilde{\sigma}_\infty \gamma$$

and

$$\tilde{\omega}_j = \tilde{\sigma}_j^*(\tilde{\omega}) = (g^{-1} \circ \sigma_j)^*(g^*(\omega)) = \sigma_j^* g^{-1*} g^*(\omega) = \sigma_j^*(\omega) = \omega_j. \quad \square$$

The curve  $t \mapsto \sigma_j(t, \theta)$  is horizontal, hence  $\omega_0(0) = \omega_\infty(\infty) = 0$  and  $\omega_j = f_j d\theta$ . As  $f_0 = \gamma^{-1} f_\infty \gamma + \gamma^{-1} \gamma'$ , we can put

$$(1.5) \quad f = f_0 - \cos^2\left(\frac{\varphi}{2}\right) \gamma^{-1} \gamma' = \gamma^{-1} \left( f_\infty + \sin^2\left(\frac{\varphi}{2}\right) \gamma' \gamma^{-1} \right) \gamma,$$

and obtain a well-defined map  $f: S^2 \rightarrow \mathfrak{g}$ .

As  $\omega_0$  has zero  $d\phi$  component,  $[\omega_0, \omega_0] = 0$ , so the curvature  $R^\omega$  has on  $U_0$  the local expression

$$R_0^\omega = d\omega_0 = \frac{\partial f_0}{\partial \varphi} d\varphi \wedge d\theta$$

and

$$\frac{\partial f_0}{\partial \varphi} = \frac{\partial f}{\partial \varphi} - \cos\left(\frac{\varphi}{2}\right) \sin\left(\frac{\varphi}{2}\right) \gamma^{-1} \gamma' = \frac{\partial f}{\partial \varphi} - \frac{1}{2} \sin(\varphi) \gamma^{-1} \gamma'.$$

The Yang-Mills functional of the connection  $\omega$  is given by

$$\mathcal{YM}(\omega) = \frac{1}{2} \int_{S^2} \|R^\omega\|^2 = \frac{1}{2} \int_{S^2} \langle *R^\omega, *R^\omega \rangle d\Omega,$$

where  $*$  is the Hodge star operator,  $d\Omega = \sin(\varphi) d\varphi \wedge d\theta$  is the standard volume form on  $S^2$  and the bracket is the inner product on the Lie algebra  $\mathfrak{g}$ . We have  $*(d\varphi \wedge d\theta) = \frac{1}{\sin(\varphi)}$ , so

$$\begin{aligned} \mathcal{YM}(\omega) &= \frac{1}{2} \int_0^{2\pi} \int_0^\pi \left\| \frac{\partial f_0}{\partial \varphi} \frac{1}{\sin(\varphi)} \right\|^2 \sin(\varphi) d\varphi d\theta \\ &= \frac{1}{2} \int_0^{2\pi} \int_0^\pi \left\langle \frac{\partial f}{\partial \varphi} - \frac{1}{2} \sin(\varphi) \gamma^{-1} \gamma', \frac{\partial f}{\partial \varphi} - \frac{1}{2} \sin(\varphi) \gamma^{-1} \gamma' \right\rangle \frac{1}{\sin(\varphi)} d\varphi d\theta \\ &= \frac{1}{2} \int_0^{2\pi} \int_0^\pi \left\| \frac{1}{\sin(\varphi)} \frac{\partial f}{\partial \varphi} \right\|^2 \sin(\varphi) d\varphi d\theta \\ &\quad - \frac{1}{2} \int_0^{2\pi} \int_0^\pi \left\langle \frac{\partial f}{\partial \varphi}, \gamma^{-1} \gamma' \right\rangle d\varphi d\theta \\ &\quad + \frac{1}{8} \int_0^{2\pi} \int_0^\pi \|\gamma' \gamma^{-1}\|^2 \sin(\varphi) d\varphi d\theta. \end{aligned}$$

We have

$$\int_0^{2\pi} \int_0^\pi \left\langle \frac{\partial f}{\partial \varphi}, \gamma^{-1} \gamma' \right\rangle d\varphi d\theta = \int_0^{2\pi} \left[ \langle f, \gamma^{-1} \gamma' \rangle \right]_0^\pi d\theta = 0$$

and

$$\begin{aligned} \int_0^{2\pi} \int_0^\pi \|\gamma^{-1}\gamma'\|^2 \sin(\varphi) d\varphi d\theta &= \int_0^\pi \sin(\varphi) d\varphi \int_0^{2\pi} \|\gamma^{-1}\gamma'\|^2 d\theta \\ &= 2 \int_0^{2\pi} \|\gamma^{-1}\gamma'\|^2 d\theta \\ &= 8\pi\mathcal{E}(\gamma), \end{aligned}$$

where  $\mathcal{E}(\gamma)$  is the *energy* of the loop  $\gamma$ . Finally  $\int_0^{2\pi} \int_0^\pi \left\| \frac{1}{\sin(\varphi)} \frac{\partial f}{\partial \varphi} \right\|^2 \sin(\varphi) d\varphi d\theta$  is a norm  $\|\cdot\|$  on the vector space  $\{f \in C^\infty(S^2, \mathfrak{g}) \mid f(0) = f(\infty) = 0\}$ . If we let  $\mathcal{H}$  be the completion, then we have proved

(1.6) **THEOREM.** *Let  $\mathcal{C}$  be the space of connections,  $\mathcal{G}_0$  the group of based gauge transformations and  $\mathcal{B} = \mathcal{C}/\mathcal{G}_0$  the orbit space, then the map  $\omega \mapsto (\gamma, f)$  induces an injection*

$$\mathcal{B} \hookrightarrow \Omega G \times \mathcal{H}$$

and

$$\mathcal{YM}(\omega) = \frac{1}{2}\|f\|^2 + \pi\mathcal{E}(\gamma).$$

The image of  $\mathcal{B}$  is not a product of  $\Omega G$  and a subspace of  $\mathcal{H}$ , but consists of pairs  $(\gamma, f)$  such that  $f_0$  and  $f_\infty$  defined by (1.5) are smooth on  $U_0$  and  $U_\infty$  respectively. Especially,  $f$  is smooth outside  $\{0, \infty\}$ .

We could replace (1.5) with

$$(1.7) \quad \tilde{\omega} = \omega_0 - h(\varphi)\gamma^{-1}\gamma' d\theta = \gamma^{-1}(\omega_\infty + h(\varphi)\gamma'\gamma^{-1})\gamma,$$

where  $h(\varphi)$  is a smooth function, which is 0 in a neighbourhood of  $\pi$  and 1 in a neighbourhood of 0. Then  $\tilde{\omega}$  is a smooth one form on  $S^2$  with values in  $\mathfrak{g}$  and with  $\tilde{\omega}(\frac{\partial}{\partial \varphi}) = 0$  at all points. Conversely, given  $\gamma \in \Omega G$  and such a one form  $\tilde{\omega}$ , then (1.7) gives a pair  $\omega_0, \omega_\infty$ , which in turn defines a connection in the bundle defined by  $\gamma$ . I.e. we have

(1.8) **THEOREM.** *The map  $\mathcal{B} \rightarrow \Omega G$  is a trivial vector bundle.*

But we do not have the nice formula for the Yang-Mills functional, because if we try to use (1.7) in the calculation, then the mixed term  $\int \langle d\tilde{\omega}, h'\gamma^{-1}\gamma' \rangle$  does not vanish.

## 2. Holomorphic trivialization

If  $G_{\mathbb{C}}$  is the complexification of  $G$  and  $LG_{\mathbb{C}}$  is the space of free smooth loops in  $G_{\mathbb{C}}$ , then  $\Omega G = LG_{\mathbb{C}}/L^+G_{\mathbb{C}}$ , where  $L^+G_{\mathbb{C}}$  is the space of loops which is the

boundary value of a holomorphic map  $D \rightarrow G_{\mathbb{C}}$  and  $D$  is the unit disk in  $\mathbb{C}$ , see [2]. We can use this description of  $\Omega G$  to get a different map  $\mathcal{C} \rightarrow \Omega G$ .

Using stereographic projection, we identify  $S^2$  with  $\mathbb{C} \cup \{\infty\}$ . Then  $U_0$  corresponds to  $\mathbb{C}$ ,  $D_0$  corresponds to the unit disk, and the meaning of 0 and  $\infty$  are unchanged.

The connection  $\omega$  on  $P$  induces a holomorphic structure on the complexified bundle  $P_{\mathbb{C}} = P \times_G G_{\mathbb{C}}$ , see [1], and the base point  $p_{\infty}$  in  $P$  is also a base point in  $P_{\mathbb{C}}$ . We still let  $\sigma_{\infty}$  be obtained by radial trivialization from the North Pole, but now  $\sigma_0$  shall be a holomorphic section over  $\overline{D}_0$  (it means that  $\sigma_0$  is smooth on  $\overline{D}$  and holomorphic on  $D_0$ ). As before  $\gamma \in LG_{\mathbb{C}}$  is defined by  $\sigma_0 = \sigma_{\infty}\gamma$ , and clearly  $\gamma$  is well-defined up to multiplication with an element of  $L^+G_{\mathbb{C}}$ , i.e. we have a well-defined element  $[\gamma] \in LG_{\mathbb{C}}/L^+G_{\mathbb{C}} = \Omega G$ .

Let  $\hat{\sigma}_0$  denote the section over  $D_0$  and  $\hat{\gamma}$  the loop in  $G$  obtained by radial trivialization. We want to find a holomorphic section of  $P_{\mathbb{C}}$  over  $D_0$ . Such a section has the form  $\hat{\sigma}_0 g$ , where  $g: \overline{D}_0 \rightarrow G_{\mathbb{C}}$ , and then  $\gamma = \hat{\gamma}g$ .

In the trivialization defined by the section  $\hat{\sigma}_0$ , the covariant derivative induced by  $\omega$  is  $\nabla = d + f_0 d\theta$ . We have  $z = \cot(\frac{\varphi}{2})e^{i\theta}$ , so

$$dz = d\left(\cot\left(\frac{\varphi}{2}\right)\right)e^{i\theta} + \cot\left(\frac{\varphi}{2}\right)ie^{i\theta}d\theta,$$

and

$$d\theta = \frac{i}{2\cot\left(\frac{\varphi}{2}\right)}(e^{i\theta}d\bar{z} - e^{-i\theta}dz) = \frac{i}{2|z|^2}(z d\bar{z} - \bar{z} dz).$$

If  $g: \overline{D}_0 \rightarrow G_{\mathbb{C}}$ , then  $\sigma_0 g: \overline{D}_0 \rightarrow P_{\mathbb{C}}$  is holomorphic if and only if

$$(*) \quad g^{-1}\frac{\partial g}{\partial \bar{z}} + \frac{iz}{2|z|^2}g^{-1}f_0g = 0.$$

In general we have  $[\gamma] \neq [\hat{\gamma}]$  because equality holds if and only if  $g|_{S^1} \in L^+G_{\mathbb{C}}$  and this need not be the case. Choose for example  $\gamma = 1$  and  $f(z) = 4i|z|^2\bar{z}A$  on  $D_0$  with  $A \in \mathfrak{g}$ , then (\*) becomes

$$g^{-1}\frac{\partial g}{\partial \bar{z}} - 2z\bar{z}g^{-1}Ag = 0.$$

A solution is  $g(z) = \exp(z\bar{z}^2A)$ , but  $g(e^{i\theta}) = \exp(e^{-i\theta}A)$  so  $g|_{S^1} \notin L^+G_{\mathbb{C}}$ .

But, if  $\omega$  is an instanton, then  $f = 0$  and  $\gamma = \exp(\theta A)$  for some  $A \in \mathfrak{g}$ , hence

$$f_0 = f + \cos^2\left(\frac{\varphi}{2}\right)\gamma^{-1}\gamma' = \frac{|z|^2}{1+|z|^2}A.$$

So (\*) becomes

$$g^{-1}\frac{\partial g}{\partial \bar{z}} + \frac{i}{2}\frac{z}{1+|z|^2}g^{-1}Ag = 0,$$

and  $g(z) = \exp(-\frac{i}{2}\log(1+|z|^2)A)$  is a solution. As  $g(e^{i\theta}) = \exp\left(-\frac{i\log 2}{2}A\right)$ ,  $g \in L^+G_{\mathbb{C}}$  and  $[\gamma] = [\hat{\gamma}g] = [\hat{\gamma}]$ .

### 3. The General Case

The process described in §2 makes sense on any closed Riemann surface  $X$ . Choose a point  $x_\infty$  on  $X$  and a local parameter around  $x_\infty$ . We shall write the local parameter as  $z^{-1}$ , thus  $z$  is a holomorphic map from a neighbourhood of  $x_\infty$  to a neighbourhood of  $\infty$  in the Riemann sphere. We may assume that  $z(x_\infty) = \infty$ , and that  $z$  is an isomorphism between a neighbourhood  $U_\infty$  of  $x_\infty$  and the region  $|z| > \frac{1}{2}$  on the Riemann sphere. The standard circle  $S^1$  can then be identified with the circle  $|z| = 1$  around  $x_\infty$  on  $X$ . We denote the part of  $X$  where  $|z| > 1$  by  $X_\infty$ , and the complement of the region where  $|z| \geq 1$  by  $X_0$ . Thus

$$\overline{X_0} \cap \overline{X_\infty} = S^1.$$

The space of smooth loops  $\gamma: S^1 \rightarrow G_{\mathbb{C}}$ , which is the boundary value of a holomorphic map  $X_0 \rightarrow G_{\mathbb{C}}$ , is denoted  $L_X^+ G_{\mathbb{C}}$ . Both  $\Omega G = LG_{\mathbb{C}}/L^+ G_{\mathbb{C}}$  and the quotient  $Gr^X = LG_{\mathbb{C}}/L_X^+ G_{\mathbb{C}}$  are subvarieties of an infinite dimensional Grassmannian  $Gr$ , see [2].

Let  $P$  be a principal  $G$ -bundle over  $X$  with a connection  $\omega$  and choose a base point  $p_\infty$  in the fiber over  $x_\infty$ . Let  $\theta \in [0, 2\pi]$  and  $r \in (\frac{1}{2}, \infty]$  be polar coordinates on  $U_\infty$  and lift the curves  $t \mapsto (\frac{1}{t}, \theta)$ , with  $t \in [0, 1]$  and  $\theta \in [0, 2\pi]$ , to horizontal curves in  $P$  starting at  $p_\infty$ . As for  $S^2$ , we get a smooth section  $\sigma_\infty$  in  $P$  over  $\overline{X_\infty}$ .

Let  $\sigma_0$  be a holomorphic section in the complexified bundle  $P_{\mathbb{C}}$  over  $\overline{X_0}$ . The transition function  $\gamma: S^1 \rightarrow G_{\mathbb{C}}$  is defined by

$$(3.1) \quad \sigma_0|_{S^1} = \sigma_\infty|_{S^1} \gamma.$$

As the section  $\sigma_0$  is only determined up to multiplication by a holomorphic map  $\overline{X_0} \rightarrow G_{\mathbb{C}}$ , the loop  $\gamma$  is only determined up to multiplication by an element of  $L_X^+ G_{\mathbb{C}}$ , but we have a well-defined element  $[\gamma] \in LG_{\mathbb{C}}/L_X^+ G_{\mathbb{C}} = Gr^X$ . We put

$$\omega_j = \sigma_j^*(\omega)$$

Not any triple  $(\gamma, \omega_0, \omega_\infty)$  can be obtained this way. At least we must have  $\omega_0 \in C^\infty(\wedge^{1,0}(\overline{X_0}, \mathfrak{g}_{\mathbb{C}}))$  and  $\omega_\infty = f_\infty d\theta$  with  $f_\infty: \overline{X_\infty} \rightarrow \mathfrak{g}$ . Given such a triple let us try to construct a  $G$ -bundle with a connection.

Using the connection  $\omega_0$  on  $\overline{X_0} \times G_{\mathbb{C}}$ , we extend  $\gamma$  by radial trivialization from  $S^1$  and obtain a smooth map  $\tilde{\gamma}: \overline{X_0} \cap U_\infty \rightarrow G_{\mathbb{C}}$ , which satisfies

$$(3.2) \quad \tilde{\gamma}^{-1} \frac{\partial \tilde{\gamma}}{\partial r} = \omega_0 \left( \frac{\partial}{\partial r} \right) \quad \text{and} \quad \tilde{\gamma}|_{S^1} = \gamma.$$

We can now use  $\tilde{\gamma}$  as the transition function in a smooth  $G_{\mathbb{C}}$ -bundle  $P_{\mathbb{C}}$ . Next we must extend  $\omega_\infty$  to a one form  $\tilde{\omega}_\infty = \tilde{f}_\infty d\theta$  with  $\tilde{f}_\infty: U_\infty \rightarrow \mathfrak{g}$ , such that

$$\omega_0 = \tilde{\gamma}^{-1} \tilde{\omega}_\infty \tilde{\gamma} + \tilde{\gamma}^{-1} d\tilde{\gamma}.$$

In view of (3.2) this is equivalent to

$$\omega_0 \left( \frac{\partial}{\partial \theta} \right) = \tilde{\gamma}^{-1} \tilde{\omega}_\infty \left( \frac{\partial}{\partial \theta} \right) \tilde{\gamma} + \tilde{\gamma}^{-1} \frac{\partial \tilde{\gamma}}{\partial \theta},$$

or

$$\tilde{f}_\infty = \tilde{\omega}_\infty \left( \frac{\partial}{\partial \theta} \right) = \tilde{\gamma} \omega_0 \left( \frac{\partial}{\partial \theta} \right) \tilde{\gamma}^{-1} + \frac{\partial \tilde{\gamma}}{\partial \theta} \tilde{\gamma}^{-1}.$$

We can get such an extension if and only if we at  $S^1$  have that

$$(3.3) \quad f_\infty = \omega_\infty \left( \frac{\partial}{\partial \theta} \right) = \tilde{\gamma} \omega_0 \left( \frac{\partial}{\partial \theta} \right) \tilde{\gamma}^{-1} + \frac{\partial \tilde{\gamma}}{\partial \theta} \tilde{\gamma}^{-1} \text{ to all orders.}$$

If (3.3) is satisfied, then we get a connection on  $P_{\mathbb{C}}$ . Finally, as  $\omega_\infty$  has values in  $\mathfrak{g}$  we have a  $G$ -structure over  $\bar{X}_\infty$ . And over  $\bar{X}_0$  the complex structure together with the connection give a  $G$ -structure. So all in all we get a  $G$ -bundle with a connection.

Furthermore, given  $(\gamma, \omega_0)$  such that the right-hand side of (3.3) lies in  $\mathfrak{g}$  at  $S^1$ , then by [3] we can find  $\omega_\infty$  satisfying (3.3), and  $\omega_\infty$  can be chosen such that it depends continuously on  $(\gamma, \omega_0)$ . Let

$$\Omega_{d\theta}^1(\bar{X}_\infty, \mathfrak{g}) = \left\{ \alpha \in \Omega^1(\bar{X}_\infty, \mathfrak{g}) \mid \alpha \left( \frac{\partial}{\partial r} \right) = 0 \right\},$$

$$\mathcal{K} = \{ (\gamma, \omega_0, \omega_\infty) \in LG_{\mathbb{C}} \times \Omega^{1,0}(\bar{X}_0, \mathfrak{g}_{\mathbb{C}}) \times \Omega_{d\theta}^1(\bar{X}_\infty, \mathfrak{g}) \mid (3.3) \text{ is satisfied} \}$$

and

$$\mathcal{L} = \left\{ (\gamma, \omega_0) \in LG_{\mathbb{C}} \times \Omega^{1,0}(\bar{X}_0, \mathfrak{g}_{\mathbb{C}}) \mid \tilde{\gamma} \omega_0 \left( \frac{\partial}{\partial \theta} \right) \tilde{\gamma}^{-1} + \frac{\partial \tilde{\gamma}}{\partial \theta} \tilde{\gamma}^{-1} \in \mathfrak{g} \text{ on } S^1 \right\}.$$

We let  $\mathcal{C}$  denote the space of connections,  $\mathcal{G}_0$  the group of based gauge transformations and  $\mathcal{B}$  the orbit space  $\mathcal{C}/\mathcal{G}_0$ .

The group  $L_X^+ G_{\mathbb{C}}$  acts on  $\mathcal{L}$  (and  $\mathcal{K}$ ) by  $(\gamma, \omega_0)g = (\gamma g, g^{-1} \omega_0 g + g^{-1} dg)$ , and we have well defined maps

$$\begin{array}{ccc} \mathcal{B} & \xrightarrow{\cong} & \mathcal{K}/L_X^+ G_{\mathbb{C}} \\ & & \downarrow \\ & & \mathcal{L}/L_X^+ G_{\mathbb{C}} \\ & & \downarrow \\ & & Gr^X. \end{array}$$

The map  $\mathcal{K}/L_X^+ G_{\mathbb{C}} \rightarrow \mathcal{L}/L_X^+ G_{\mathbb{C}}$  is an affine bundle, and the vector space of translations is in each fiber

$$\{ \omega_\infty \in \Omega_{d\theta}^1(\bar{X}_\infty, \mathfrak{g}) \mid \omega_\infty \text{ vanishes to all orders on } S^1 \}.$$

Likewise  $\mathcal{L}/L_X^+ G_{\mathbb{C}} \rightarrow Gr^X$  is an affine bundle with fiber

$$\{ \omega_0 \in \Omega^{1,0}(\bar{X}_0, \mathfrak{g}_{\mathbb{C}}) \mid \omega_0 \in \gamma^{-1} \mathfrak{g} \gamma - \gamma^{-1} \gamma' \text{ on } S^1 \}.$$

Hence we have



(3.4) THEOREM. *The map  $\mathcal{B} \rightarrow Gr^X$  is a homotopy equivalence.*

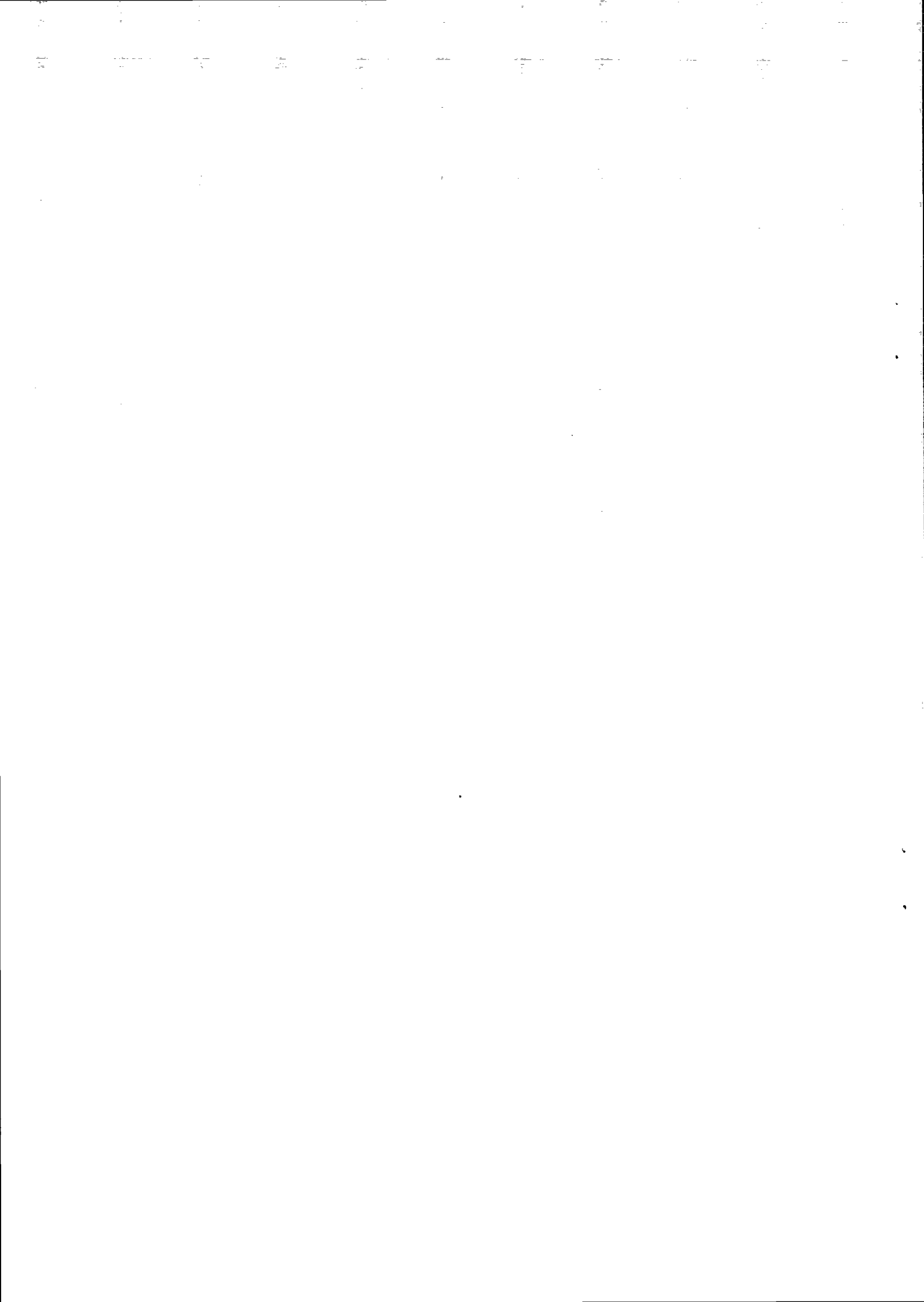
It is already known that  $\text{Map}(X, BG)$  is homotopy equivalent to both  $\mathcal{B}$ , see [1], and  $Gr^X$ , see [2]. We have now shown the third homotopy equivalence directly.

It would be nice if one could get a theorem like (1.6) for a general Riemann surface. In §2 we saw that the map we get by radial trivialization is different from the map given by holomorphic trivialization, so perhaps theorem (1.6) is too much to ask for. But the two maps agree on the moduli space of instantons, so we could look for a function  $\mathcal{F}: Gr^X \rightarrow \mathbb{R}$ , such that the moduli space of instantons is mapped bijectively to the critical points of  $\mathcal{F}$ .

The energy-function on  $\Omega G$  is the restriction of a function defined on all of  $Gr$ , see [2], so the restriction to  $Gr^X$  could be a candidate for  $\mathcal{F}$ .

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